## ON THE NATURE OF THE DUST IN THE DEBRIS DISK AROUND HD 69830

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## ABSTRACT

We have used the infrared mineralogical model derived from the *Spitzer* IRS observations of the Deep Impact experiment to study the nature of the dust in the debris found around the K0 V star HD 69830. Using a robust approach to determine the bulk average mineralogical composition of the dust, we show it to be substantially different from that found for comets 9P/Tempel 1 and C/Hale-Bopp 1995 O1 or for the comet-dominated YSO HD 100546. Lacking in carbonaceous and ferrous materials but including small icy grains, the composition of the HD 69830 dust most closely resembles that of a disrupted P- or D-type asteroid. The amount of mass responsible for the observed emission is the equivalent of a 30 km radius, 2500 kg m<sup>-3</sup> sphere, while the radiative temperature of the dust implies that the bulk of the observed material is at ~1.0 AU from the central source, coincident with the 2:1 and 5:2 mean motion resonances of the outermost of three Neptune-sized planets detected by Lovis and coworkers. In our solar system, P- and D-type asteroids are both large and numerous in the outer main belt and near Jupiter (e.g., the Hildas and Trojans) and have undergone major disruptive events to produce debris disk–like structures (cf. the Karin and Veritas families 5–8 Myr ago). The short-lived nature of the small and icy dust implies that the disruption occurred within the last year, or that replenishment due to ongoing collisional fragmentation is occurring.

Subject headings: comets: general - infrared: general - minor planets, asteroids - techniques: spectroscopic

## 1. INTRODUCTION

One of the most dramatic Spitzer discoveries to date has been the extreme level of zodiacal emission around the nearby K0 V star HD 69830 (Fig. 1). The dust cloud around HD 69830 has more than 1000 times the emission of our own zodiacal cloud and shows a plethora of solid state features attributable to small, hot, crystalline silicate grains located within 0.5-1.0 AU of the parent star (Beichman et al. 2005; Fig. 2). Such intense emission from dust in the inner solar system is exceedingly rare (Bryden et al. 2006; Beichman et al. 2006b). Interest in this cloud and its link to the evolution of planetary systems was heightened recently by the announcement that three Neptune mass planets orbit within 0.63 AU of the star (Lovis et al. 2006). With an estimated age of 3–10 Gyr, a near-solar primary ( $T_* \sim 5385$  K,  $L_* = 0.60 L_{\odot}$ ,  $M_* \sim$  $0.86 M_{\odot}$ ,  $R_* \sim 0.89 R_{\odot}$ ), and a near-solar metallicity ([Fe/ H] = 0.015, Beichman et al. 2005;  $[Fe/H] = -0.05 \pm 0.02$ , Lovis et al. 2006), the HD 69830 system promises to be an excellent laboratory for the study of the interaction been a debris disk and a solar-type planetary system. One of the major unanswered questions about this system is whether the dust seen by Spitzer comes from collisions within a particularly massive asteroid belt or from a swarm of comets released by planet-disk interactions, and whether it is transient or persistent in nature. Here we use the infrared mineralogical model derived from the Spitzer Infrared Spectrograph (IRS) observations of the Deep Impact experiment (A'Hearn et al. 2005; Lisse et al. 2006) to investigate the nature of the dust found around HD 69830.

# 2. DATA

This paper uses the low-resolution ( $R \sim 100$ ) Spitzer IRS  $(7-35 \,\mu\text{m})$  measurements previously described in Beichman et al. (2005) (Figs. 1 and 2). The reader is referred to that paper and to Beichman et al. (2006a) for details of the data reduction. The separation of the emission due to the stellar photosphere from the "excess" due to the orbiting debris disk is estimated to be accurate to 1%-2% of the photospheric level (Fig. 1). As a test of the data reduction, we rereduced the IRS data using the S13 pipeline from the Spitzer Science Center, which uses a different flat-fielding technique. After removing the stellar photosphere, we compared the two different versions of the spectrum of the excess. The agreement was generally excellent: between the two spectra there are only 3 out of 270 data points that deviated by more than 3  $\sigma$  from one another and only 15 points that deviated by more than  $2\sigma$ . The resulting spectral models are only different, within the errors of the fit, for the amorphous carbon smooth component, which mimics the continuum at short wavelengths.) In the modeling discussed below we use the original published spectrum (Fig. 2). The quoted uncertainties in the derived composition information (Table 1) are consistent with either reduced spectrum and the associated errors.

## 3. MODELING

Beichman et al. (2005) modeled the emission spectrum of the material in the debris disk by assuming a zeroth-order model using dust with comet Hale-Bopp's optical parameters and 400 K temperature (Fig. 2). From this, they deduced a total mass for the debris disk that exceeded the mass of any known comet, suggesting an asteroidal rather than a cometary origin for the debris. However, the quality of their fit was less than ideal, as the nature of the dust in the two systems is significantly different on detailed examination. In this work, a more sophisticated spectral modeling approach was taken, motivated by the successful compositional analysis of the material excavated by the Deep Impact experiment from the interior of comet 9P/Tempel 1 (Lisse et al.

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FIG. 1.—SED of HD 69830 derived from *Hipparcos* measurements (optical photometric points), 2MASS (*J*, *H*, and  $K_s$  photometric points), *Spitzer* 7–35  $\mu$ m spectroscopy (*orange curve*), *IRAS* 12 and 25  $\mu$ m photometry (Mannings & Barlow 1998), and *Spitzer* 24 and 70  $\mu$ m photometry. Error bars are 1  $\sigma$  and are smaller than the symbol size for some points. All *Spitzer* data are from Beichman et al. (2005). Overlaid is the Kurucz model (*blue*) fit to the short-wavelength photometry used to remove the stellar photospheric contribution. The 2MASS data are very uncertain because of the brightness of the star.

2006) and comet Hale-Bopp and young stellar object (YSO) HD 100546 (Lisse et al. 2007). As done for Tempel 1 and Hale-Bopp, the HD 69830 excess flux was calculated by first removing the photospheric contribution using a Kurucz model fit and then converting the remaining excess flux to an emissivity spectrum by dividing the measured fluxes by a best-fit blackbody (here found to be 400 K from  $\chi^2$  minimization testing, in good agreement with the findings of Beichman et al. 2005). We work in emissivity space because it is very easy to pick out the contributions of the different mineral species once the gross effects of temperature are removed.

The resulting emissivity spectrum (Fig. 3) was then compared to a linear combination of emission spectra of over 80 candidate mineral species (there are no obvious gas emission lines), selected for their reported presence in YSOs, solar system bodies, dusty disk systems, and interplanetary dust particles. The results of the modeling are returned as the coefficients (weights) required to multiply each normalized mineral species emissivity spectrum in the best-fit linear sum spectrum (Fig. 3 and Table 1). Assuming crystalline densities for the mainly small (~1  $\mu$ m or less radius) dust, we can convert these weights, which are the observed surface area of each species referenced to a 400 K blackbody at the distance of HD 69830, into the relative number of moles of each species. To derive a lower limit for the total mass of each species, we integrate the mass represented by the best-fit particle size distribution (PSD) from 0.1 to 10  $\mu$ m, the range of dust particle sizes the Spitzer spectrum is sensitive to, determining the absolute normalization by the amount of material required to produce the observed absolute flux at 12.6 pc distance.

When doing this comparison, the detailed properties of the emitting dust, i.e., the particle composition, size distribution, temperature, and porosity, all had to be addressed, as they can strongly affect the observed infrared flux. These quantities are thus products of the model solutions. (The reader is referred to the discussion given at length in Lisse et al. [2006] and in the online supplemental material [SOM] for additional details of the calculations.) Verification and validation of the model, as well as a measure of the range of modeling compositions found for cometary systems, can be found in the comparison of comets Hale-Bopp, Tempel 1, and the YSO HD 100546 (Lisse et al. 2007).

We present here a model (Figs. 3 and 4) consisting of the fewest and simplest dust species possible that result in a consistent fit to the observational data. A set of components was tested exhaustively before the addition of a new species was allowed, and only species that reduced the  $\chi^2_{\nu}$  below the 95% confidence



FIG. 2.—Comparison of the HD 69830 and Hale-Bopp emission spectra. *Black*: Spectrum of the excess of HD 69830 according to Beichman et al. (2005). *Red*: For comparison, the spectrum of Comet C/1995 O1 (Hale-Bopp), following Crovisier et al. (1997). *Blue*: Hale-Bopp spectrum normalized to 550 K to best match the HD 69830 spectrum. Note the comparative excess of emission for Hale-Bopp at  $8-10 \mu m$  due to pyroxene and at  $12-15 \mu m$  due to water ice.

TABLE 1										
COMPOSITION	OF THE	Best-Fit	Model	то тни	Spitzer	IRS	HD	69830	Spectrum	[

Species	Weighted <sup>a</sup> Surface Area	Density (g cm <sup>-3</sup> )	Molecular Weight	$N_{\rm moles}^{\ \ b}$ (Relative)	Model T <sub>max</sub> (K)	Model $\chi_{\nu}^{2c}$ if Not Included
		Detectio	ons			
Olivines:						
Amorph olivine (MgFeSiO <sub>4</sub> )	0.19	3.6	172	0.40	370	13.3
ForsteriteKoike (Mg <sub>2</sub> SiO <sub>4</sub> )	0.20	3.2	140	0.45	370	4.83
Forsterite038 (Mg <sub>2</sub> SiO <sub>4</sub> ) <sup>d</sup>	0.33	3.2	140	0.74	370	15.9
Fayalite (Fe <sub>2</sub> SiO <sub>4</sub> )	0.04	4.3	204	0.07	370	1.29
Pyroxenes:						
Bronzite $(Mg_{1-x}Fe_xSi_2O_6)^d$	0.13	3.5	232	0.19	340	4.47
FerroSilite (Fe <sub>2</sub> Si <sub>2</sub> O <sub>6</sub> )	0.09	4.0	264	0.14	340	2.46
Diopside (CaMgSi <sub>2</sub> O <sub>6</sub> )	0.05	3.3	216	0.075	340	1.53
Carbonates:						
Magnesite (MgCO <sub>3</sub> )	0.07	3.1	84	0.24	370	1.54
Dolomite (CaMgC <sub>2</sub> O <sub>6</sub> ) <sup>d</sup>	0.04	2.9	184	0.064	370	1.25
Water ice (H <sub>2</sub> O)	0.08	1.0	18	0.44	220	2.28
Amorphous carbon (C)	0.07	2.5	12	≤1.5	410	1.78
		Upper Li	mits			
Sulfides (Mg <sub>50</sub> Fe <sub>50</sub> S)	0.03	4.5	72	$\leq 0.02$	370	1.22
PAH (C <sub>10</sub> H <sub>14</sub> ), ionized	0.05	1.0	$\langle 178 \rangle$	$\leq 0.028$		1.10

<sup>a</sup> Weight of the emissivity spectrum of each dust species required to match the HD 69830 emissivity spectrum.

<sup>b</sup>  $N_{\text{moles}}(i) \sim \text{density}(i)/\text{molecular weight}(i) \times \text{normalized surface area}(i)$ . Errors are  $\pm 10\%$  (2  $\sigma$ ).

<sup>c</sup> Total best-fit model  $\chi^2_{\nu} = 1.09$ .

<sup>d</sup> Not found in cometary systems to date.

limit (CL) were kept. It is important to note that we have not necessarily found the exact mineral species with our technique; rather, we claim to have found the important mineral classes present in the dust and the gross atomic abundance ratios. It is also important to emphasize that while the number of parameters (composition, temperature, PSD) may seem large, instead there are very few detected species for the 270 independent spectral points and 13 strong features obtained at high signal-to-noise ratio by the *Spitzer* IRS over the 7–35  $\mu$ m range. It was in fact extremely difficult to fit the observed spectrum within the 95% CL of  $\chi^2_{\mu} = 1.09$ .

#### 4. DERIVED PROPERTIES OF THE DUST

#### 4.1. Temperature

We modeled the dust excess around HD 69830 as both an extended disk of dust and a localized dust torus. The toroidal models were motivated by the single-temperature distributions found for many of the objects studied by Beichman et al. (2006a) and Chen et al. (2006), as well as the narrow dust structures found in many of the Hubble Space Telescope images of debris disks (Kalas et al. 2005, 2006). We find that only narrow toroidal models fit the data within the 95% CL. Our best-fit temperature of the smallest dust particles  $(0.1-1.0 \,\mu\text{m})$ , which superheat substantially above local thermal equilibrium (LTE), is  $T_{\text{max}} = 370 \text{ K}$ for the olivines, 340 K for the pyroxenes, and 410 K for the amorphous carbon. The largest particles in our calculation, 1000  $\mu$ m, are set to  $T_{\text{LTE}}$ , and dust of intermediate size is scaled between the two extremes (Lisse et al. 2006). Using the Tempel 1 ejecta temperatures as a guide (where we found olivines and pyroxenes at  $\sim$ 340 K and amorphous carbon at 390 K at 1.51 AU from the Sun, where LTE = 230 K), and allowing for an HD 69830 stellar radius that is 0.89 the Sun's, a luminosity that is 0.60  $L_{\odot}$ , and a  $T_{\text{eff}}$  that is 5385 K (Lovis et al. 2006), we estimate a location for the dust at 0.93–1.16 AU from HD 69830. This range of distances is coincident with the strong 2:1 and 5:2 mean motion resonances of HD 69830d predicted by Lovis et al. (2006) from their Monte Carlo modeling of stable dust particle orbits and the inner edge of a large distribution of low-eccentricity stable orbits (Fig. 5).

Interestingly, the best-fit single continuum temperature to the 7–35  $\mu$ m *Spitzer* spectrum is 400 K, close to the amorphous carbon temperature. As seen in the Tempel 1 ejecta from the Deep Impact experiment (Lisse et al. 2006), the amorphous carbon dominates the continuum behavior as it is the hottest material and the most continuum-like, thus contributing most to the shortwavelength emission. It is the short-wavelength emission that is most diagnostic of an SED's continuum temperature because of the steep falloff of the blackbody on the Wien law side of the emission peak and the relatively shallow and temperature-independent behavior on the long-wavelength Rayleigh-Jeans side of the peak.

For scale, putting the dust at the mean distance of the outermost planet, HD 69830d, at  $\sim 0.63$  AU from the central source (Fig. 5), and assuming graybody behavior for the dust (i.e., ignoring any superheating due to differences in the particle optical and IR emissivities) yields dust at  $T_{\text{LTE}} = 245/(0.63 \text{ AU})^{1/2} = 309 \text{ K}$ . Dust at the location of the middle planet HD 69830c would have  $T_{\rm LTE} =$  $245/(0.19 \text{ AU})^{1/2} = 562 \text{ K}$ ; and dust at 1 AU, the 2 : 1 mean motion resonance for the outermost Lovis planet (HD 69830d), would have  $T_{LTE} = 245$  K. Thus, had we not done the detailed compositional modeling and assumed instead that the 400 K continuum temperature was equal to  $T_{LTE}$ , we would have placed the dust erroneously at a distance of approximately 0.4 AU from the central source. The presence of abundant water ice, however, in our best-fit compositional solution is a clue that this would be incorrect, as water ice would be stable for only seconds at 0.4 AU.



FIG. 3.—Emissivity spectrum for HD 69830 excess, as measured by the *Spitzer* IRS. The data and model shown are for the smallest photospheric subtraction, demonstrating the upper limits to the abundance of PAHs, amorphous carbon, and sulfides in the system. Error bars are  $\pm 2 \sigma$ . *Black: Spitzer* excess spectrum, divided by a 400 K blackbody. *Orange dashed line:* Best-fit model spectrum. *Colored lines:* Emission spectra for the constituent species, scaled by the ratio  $B_2[T_{dust}(a)_i]/B_{\lambda}(T_{continuum})$ , with  $T_{continuum} = 400$  K. The best-fit model individual species curves have been scaled by a factor of 2 vs. the data for emphasis. *Purple:* Amorphous silicates of pyroxene or olivine composition. *Light blue:* Crystalline pyroxenes: ferrosilite, diopside, and orthoenstatite, in order of 20  $\mu$ m amplitude. *Dark blue:* Crystalline olivines, forsterite, and fayalite, in order of 20  $\mu$ m amplitude. *Red:* Amorphous carbon. *Dark orange:* Water ice. *Light orange:* Water gas. *Yellow:* PAHs. *Bright green:* Carbonates: siderite and magnesite, by order of 7  $\mu$ m emissivity amplitude. *Olive green:* Sulfides.

Extrapolating the best-fit model spectrum longward of 35  $\mu$ m using an assumed absorption efficiency  $Q_{abs} = Q_{abs}(35 \ \mu m/\lambda)$ , where  $\lambda$  is the wavelength, as a lower limit and  $Q_{abs} = \text{constant}$  as an upper limit, we estimate a 70  $\mu$ m flux of 5 ± 4 mJy, while the *Spitzer* MIPS measurements of Beichman et al. (2005) yield a value of 7 ± 3 mJy for the 70  $\mu$ m excess. We thus concur with the conclusion that there is no large amount of cold (i.e., Kuiper Belt) dust contributing to the excess emission detected by *Spitzer* and that the emission is dominated by the single debris disk at  $T \sim 400$  K.

# 4.2. Size Distribution and Total Mass

The best-fit PSD of  $\sim a^{-3.9\pm0.2}$  used for all species argues for dust highly dominated by small particles in its emitting surface area and by particles of all sizes in its mass. The predominance of small particle surface area is why we see the strong emission features: the dominant particles are optically thin, and the feature-to-continuum ratio is high. This very steep PSD is unusual; more gradual ones are typically found for dusty disk systems. On the other hand, a system in collisional equilibrium would demonstrate a PSD  $\sim a^{-3.5}$  (Dohnanyi 1969; Durda & Dermott 1997); for "real" systems a size distribution even steeper than  $a^{-3.5}$  at small sizes is expected in a collisional cascade both because of its truncation at the blowout limit (Thebault et al. 2003) and because of the dependence of particle strength on size (O'Brien & Greenberg 2003).

Integrating the mass implied by the best-fit PSD, we find an estimated dust mass of  $\sim 3 \times 10^{17}$  kg definitively detected by the *Spitzer* IRS spectrum (i.e., mass in particles of  $0.1-10 \mu m$  in size that contribute appreciably to the  $\chi^2_{\nu}$  value of the fit). However, the majority of the mass is in the largest, optically thick particles if they are present  $(dm/d \log a \sim a^0)$ . While the *Spitzer* IRS measurements are not very sensitive to these large grains, the amount of mass in the larger grains ( $20 \mu m-10 m$ ; Beichman et al. 2005) can be estimated by assuming a PSD of constant slope and summing the mass of all particles up to a cutoff of 10 m. In this case we find a total disk mass of  $1.8 \times 10^{18}$  kg, or  $3 \times 10^{-7} M_{\oplus}$ . We find a total estimated  $0.1-10 \mu m$  surface area of  $1.2 \times 10^{20} m^2$ . Assuming that the dust is in a ring centered at 1.0 AU from the central star (Fig. 5), in an annulus 0.1 AU wide and with total surface



Fig. 4.—Residual spectrum found after subtraction of the best-fit silicate components, the dominant species in the emission. Scale, color, and labels are the same as in Fig. 3. The sinusoidal features at  $15-20 \ \mu m$  are artifacts ("fringes") of the IRS instrument. The true signal is near the average of the excursions.

area  $1.4 \times 10^{21}$  m<sup>2</sup>, this implies a spatial filling factor for the dust of  $0.9 \times 10^{-2}$ . In comparison, the coma of a bright, active comet, like Halley or Hale-Bopp, is observed to have an average spatial filling factor of the coma dust in the central  $10^5$  km on the order of  $10^{-7}$  (Lisse et al. 1998, 2004). It is only in the collisionally thick parts of the coma, within 10 km of the nucleus, where the coma dust spatial filling factor reaches values  $\geq 10^{-2}$ .

### 4.3. Compositional and Atomic Abundance

The derived compositional abundances from our best-fit dust model, with  $\chi^2_{\nu} = 1.09$ , are given in Table 1. As a test of the robustness of our best-fit model, the last column in Table 1 also gives the best-fitting  $\chi^2_{\nu}$  found after deleting a particular species from the compositional mix and refitting the data. Unlike the case of the comets Tempel 1 and Hale-Bopp and the comet-dominated YSO HD 100546 (Lisse et al. 2007), there is no definitive evidence for water gas, polycyclic aromatic hydrocarbons (PAHs), phyllosilcates, or sulfides (Figs. 3 and 4). The relative atomic abundance we find in the circumstellar dust is H : C : O : Si : Mg : Fe : S : Ca : Al = 0.35 : 0.90 : 4.3 : 1.0 : 1.2 : 0.4 : 0.0 : 0.06 : 0.0 (for Si = 1.0). Also unlike the comet systems, the Si : Mg : Fe : Ca : Al refractory element ratios are inconsistent with solar system abundances determined from the Sun and from C1 chondrites (Anders & Grevesse 1989), with Fe, S, and Al being highly deficient (Fig. 6).

#### 4.4. Detailed Mineralogy

Here we discuss a few of the standout materials for the case of HD 69830. We refer the reader to Figures 3 and 4 while reading this section.

### 4.4.1. Silicates

We find a best-fit model with a spectrum dominated by emission from silicates, predominantly Mg-rich olivines (80%) and crystalline pyroxenes (20%) (Fig. 3). Unlike our results on cometary systems (Lisse et al. 2006, 2007), amorphous pyroxene is conspicuously lacking. The ratio of total olivine to pyroxene abundance is  $\sim 4 : 1$ , as compared to the ratios of  $\sim 1 : 1$  in the Tempel 1 subsurface ejecta (Lisse et al. 2006),  $\sim 2 : 1$  in the dust emitted into the coma of comet C/1995 O1 (Hale-Bopp), and 1:1 in the YSO HD 100546 (Lisse et al. 2007). Both findings are direct evidence for high-temperature processing (greater than the 900 K annealing temperature of pyroxenes) of the rock-forming species.





FIG. 5.—(*a*) Schematic of the HD 69830 system, as it is understood from the present work and that of Lovis et al. (2006). In order from the center of the system: *orange-yellow center filled circle*: HD 69830a K0 V primary, at  $M_* \sim 0.86 M_{\odot}$ ,  $R_* \sim 0.89 R_{\odot}$ ,  $T_* \sim 5385$  K,  $L_* = 0.60 L_{\odot}$ ; *filled blue circles*: HD 69830b at 0.0785 AU and  $M \ge 10.2 M_{\oplus}$ , HD 69830c at 0.186 AU and  $M \ge 11.8 M_{\oplus}$ , and HD 69830d at 0.63 AU and  $M \ge 18.1 M_{\oplus}$ ; *fuzzy brown curve*: thick dust belt giving rise to the IR excess detected by *Spitzer*, located at ~1 AU, coincident with the 2:1 resonance of HD 69830d, and of total mass  $\ge 3 \times 10^{17}$  kg. Note that the planets are shown as following idealized simple circular trajectories. (*b*) Stability analysis for massless particles in the HD 69830 system, after Lovis et al. (2006). The color grid corresponds to this stability index: red denotes the most unstable orbits with a close encounter with the central star or a planet, while dark blue corresponds to very stable orbits. The particles can survive for an extended time in the stable region beyond 0.8 AU. As for the solar system main asteroid belt, several mean motion resonances with the outer-most planet can be identified that would create gaps or accumulation in a potential asteroid belt: 1:2 (0.40 AU), 2:3 (~0.48 AU), 1:1 (~0.63 AU), 3:2 (~0.82 AU), 2:1 (~1.00 AU), and 5:2 (~1.18 AU).

The relative lack of pyroxenes is what causes the deficiency in the HD 69830 spectrum versus the Hale-Bopp spectrum at 8–10  $\mu$ m (Fig. 2). The crystalline pyroxenes that are present contain a broad range of mineralogical components, including diopside, ferrosilite, and highly processed bronzite, but not simple ortho-

or clinopyroxenes that dominate in comets. Two different kinds of Mg-rich olivine were needed to accurately fit the spectrum, one of which has not been found in cometary systems, testimony to the superabundance of olivine in the dust. There is a significant fraction of amorphous olivine detected,  $\sim$ 25% of all olivine,



FIG. 6.—Atomic abundances for the circumstellar dust around HD 69830, compared to the three comet systems studied to date with the Tempel 1 ejecta compositional model. Abundances are given vs. solar, assuming Si = 1.0. Error bars for the relative measures are  $\pm 20\%$  (2  $\sigma$ ). *Diamonds*: Tempel 1; *squares*: HD 100546; *triangles*: Hale-Bopp; *asterisks*: HD 69830. The nominal solar value is denoted by the dashed line. HD 69830 is markedly different than the range of cometary abundances in Fe, S, and Al. Note that the C abundances for Hale-Bopp and HD 100546 may be artificially high due to the *ISO* SWS calibration (see Lisse et al. 2007), while Tempel 1 hydrogen and oxygen abundances include the contribution of H<sub>2</sub>O gas. The significant cometary reservoirs of gas-phase CO and CO<sub>2</sub> are not accounted for here.

arguing that some of the rock-forming elements were not raised to temperatures of 1200 K and above for significant amounts of time.

### 4.4.2. Iron and Sulfur-bearing Species

Some of the major sources of iron found in the cometary systems, the carbonate siderite (FeCO<sub>3</sub>) and the Fe sulfides (pyrrohtite, ningerite, and pentdantalite), are totally lacking in the HD 69830 debris disk. Iron is only obviously present in the minority ironbearing silicates and is clearly underabundant versus solar and C1 abundances and the other elements (Fig. 5). There is no obvious source of S in the debris disk spectrum.

#### 4.4.3. Carbon-bearing Species

The major sources of carbon in the comet systems, amorphous carbon and PAHs, are lacking in the HD 69830 best-fit model to the spectrum published in Beichman et al. (2005). The only carbon clearly present is in the form of carbonate and is clearly underabundant versus solar and C1 abundances and the other elements. Even without 5–8  $\mu$ m spectral coverage, it is clear that the 8.6, 11.3, and 12.5  $\mu$ m PAH emission lines are lacking at the spectral sensitivity and resolution of the IRS low-resolution modules. (Higher resolution observations with IRS are planned that might reveal a low level of PAH emission.)

The presence or absence of amorphous carbon is important in that except for this component, there is a good match, including the carbonates, between the best-fit compositional model and the composition of outer main belt asteroids. As the amorphous carbon signature is similar to a smooth continuum, it is possible that some of the emission from amorphous carbon has been misinterpreted as stellar photosphere emission and used to set the stellar contribution baseline. Since there is currently a low level uncertainty in the removal of the stellar photospheric contribution, which is critical to the detection of amorphous carbon, we quote here an upper limit for the amount of amorphous carbon in the debris disk by moving the photospheric model to the lowest possible allowed  $2\sigma$  limit (2%) below the nominal normalization. Doing so, we find the exact same mix of materials but with a detection of amorphous carbon. We conclude that the maximum amount of amorphous carbon possible in the system is a significant 1.8 moles (relative), or about 38% by mole fraction of the total material, 6% by surface area. Carbon at this level of abundance is very typical of that seen in the cometary bodies (Lisse et al. 2006, 2007). The minimum possible amount of amorphous carbon present is zero.

#### 4.4.4. Water Ice

The clear presence of water ice at  $11-15 \ \mu m$  is somewhat surprising at first look, given that the smallest dust particles appear to be as hot as 400 K and the average dust particle temperature is ~340 K, much above the ~200 K temperature at which water ice freely sublimates at 0 torr. However,  $T_{LTE}$  at 1 AU is 245 K, still somewhat above the ice stability temperature, but close enough to it that cooling by evaporative sublimation can stabilize the ice if it is isolated from the other hot dust. We saw the same large dichotomy in dust temperatures for the Tempel 1 ejecta in the Deep Impact experiment (Lisse et al. 2006), and in the ISO observations of YSO HD 100546 (Lisse et al. 2007).

If the water ice is unassociated with the debris disk at 1 AU but is instead present in the "Kuiper belt" for the system and is serendipitously along the line of sight, then we expect long-lived ice without much water gas formed by evaporation. However, significant icy dust in a Kuiper belt at  $\sim$ 50 AU from the star is in direct contradiction to the 3  $\sigma$  upper limit on the 70  $\mu$ m flux

observed for the system by Beichman et al. (2005). Although water ice has a strong emission feature at 65  $\mu$ m, the lack of 70  $\mu$ m emission is consistent with a simple extrapolation of emission from hot, small grains and the absence of any colder (larger) grains.

Unfortunately, we do not have  $5-8 \mu m$  measurements of the HD 69830 system in hand at this time. This region is highly diagnostic not only of water gas but also of carbonate and PAH materials found in the comets and HD 100546 (Lisse et al. 2007). Without  $5-8 \mu m$  spectroscopy, we cannot determine the presence or absence of related water gas in the system, but we hypothesize from our knowledge of water ice's thermal stability that there must be an abundance of water gas as well. On the other hand, we can say that water ice is not superabundant. We do know that the main emission from water ice is at  $11-15 \mu m$ , and the superabundance of water ice particles in the *Infrared Space Observatory (ISO)* Hale-Bopp spectrum is what caused the mismatch seen in this region in the Beichman et al. (2005) dust model for HD 69830.

## 5. RESULTS AND DISCUSSION

Taken together, the compositional and physical nature of the dust determined from the *Spitzer* spectrum provides strong evidence for differentiation and processing of material in ongoing near-collisional equilibrium. This disk is derived from material that is not primordial. Here we discuss further the implications for the putative parent body source of the dust.

## 5.1. Dynamical Considerations

The derived HD 69830 minimum dust mass of  $3 \times 10^{17}$  kg is equivalent to the mass of a very large 60 km radius comet, assuming a density of 0.35 g cm<sup>-3</sup> (from Deep Impact; A'Hearn et al. 2005; J. E. Richardson & H. J. Melosh 2007, in preparation). While larger particles may be present, we cannot prove this with the IRS measurement alone. Sensitive millimeter/submillimeter observations (with the JCMT SCUBA2, for example) would help us to understand the larger particle sizes better. The critical point here, however, is that we cannot definitively say that the mass of material is much larger than that of a comet in our solar system: the very large comet Hale-Bopp may have been of this size. If the body were an asteroid, instead, with density 2.5 g cm<sup>-3</sup>, then the minimum radius required (assuming total disruption) would be 30 km. We can say that an asteroidal body of such size is found very commonly in the solar system.

The observed dust must be relatively fresh. Due to the effects of solar radiation pressure, the lifetime for dust of radius  $\leq 1 \mu m$  (particles with  $\beta = F_{rad press}/F_{grav} > 0.5$ ; Finson & Probstein 1968; Lisse et al. 1998) in the solar system in the vicinity of 1.0 AU can be measured in weeks to months. The situation for the near-solar K0 V HD 69830 system will be little different. The lifetime for dust spiraling into the Sun due to Poynting-Robertson (PR) drag is given by  $\tau_{PR} = (10^5 \text{ yr})a(\mu m)$  (Burns et al. 1979), so even dust as large as 1 cm is cleared out within 1 Gyr. With a system age of 3-10 Gyr, the observed excess cannot be due to primordial dust created during the formation of the system. The presence of significant quantities of short-lived water ice (lifetimes of seconds to months vs. evaporation inside the HD 69830 ice line at 1.5 AU) also argues for recent generation of these small grains.

New work by Wyatt et al. (2007) also demonstrates that it is highly unlikely that an asteroid belt formed from an original nebula could create the observed excess at 3 Gyr in steady state, as collisional processing would have efficiently removed most of the mass of this belt over such a timescale. The survey of Bryden et al. (2006) shows that Sun-like stars like HD 69830, that have hot dust at a few AU or less from the primary are rare, found in just 2% of systems. Similar results were found by Rieke et al. (2005) and Song et al. (2005). It is considered that the majority of such systems are also unlikely to be explained as massive asteroid belts (Wyatt et al. 2007). Thus, studies of the aggregate population of dusty disks also argue for an extremely recent, stochastic formation of the HD 69830 dust complex.

## 5.2. Solar System Analog for the Parent Body

Given that the likelihood of serendipitously observing the breakup of an asteroid within weeks of its occurrence is unlikely, and that Mannings & Barlow (1998) found a 25  $\mu$ m excess for the system in 1983, we conclude that the small particle and icy dust components must be undergoing continual replenishment on timescales of decades or longer. The most likely physical source of continuous replenishment is a dense, localized debris belt undergoing frequent collisions between rapidly moving fragments of size 100  $\mu$ m-10 m, forming fine dust in a manner similar to the zodiacal bands in our solar system. In the case of our zodiacal bands, such collisions, occurring 5-8 Myr ago, resulted from the collisional formation of the Karin and Veritas family of asteroids from larger parent bodies (Nesvorný et al. 2006). The shepherding action of the outermost Neptune-sized planet in the system may also be important in the concentration and dynamical stirring of the debris fragments. By far the largest range of eccentric particle orbits are stably bound in the 2:1 resonance at 1 AU, where we place the dust from our modeling, than in any other resonance for the system (Lovis et al. 2006). The large range in eccentricities will cause more frequent and violent collisions for this resonance, as bodies captured into the resonance on very different orbits intersect. We thus find the collisional disruption of a body in the resonance and the localization of the collision products into the 2:1 mean motion resonance (Fig. 5) very plausible.

Given the compositional results of our modeling for the majority species constituting the dust in the debris disk around HD 69830, and using the small bodies in our solar system as a guide, we find the best match to the compositional model derived here to be a P- or D-type asteroid (in the Tholen taxonomy; e.g., Lodders & Fegley 1998, Table 13.7, p. 250, and references therein) with low carbon abundance. Comets and their immediate ancestors the Centaurs and Kuiper Belt objects are ruled out due to the lack of amorphous carbon, sulfides, and amorphous pyroxene (Zolensky et al. 2006; Lisse et al. 2006, 2007). S-type asteroids, also abundant in the inner main belt, do not contain water or products of aqueous alteration like carbonates. C-type asteroids contain much more carbon and phyllosilicates than found for the HD 69830 dust.

The P/D-type asteroids represent the most numerous asteroidal bodies in the outer main belt and in the Jovian Trojan swarms, and a typical body is of 20–100 km radius, so it is entirely plausible that the mass of debris measured for the HD 69830 system could have come from disruption of one such body, even for the larger size estimates of the original parent body determined by Beichman et al. (2005). Also, the recent finding of "comet-like asteroids" by Hsieh & Jewitt (2006) in the solar system, spectroscopic evidence for abundant hydrogen in asteroids (Rivkin et al. 2003), and the evidence for aqueous alteration on Ceres (Russell et al. 2004; McCord & Sotin 2005) make the finding of appreciable water ice produced by ongoing collisional fragmentation of a P or D asteroid credible.

The one major issue with this determination is the location of formation of the P/D-type asteroids: they represent the majority species in the outer main belt of our solar system, at 3-5 AU from the Sun today, outside the current ice line, and were presumably formed where conditions were cold enough to trap and hold most

of the volatile water present. To have the dust from such a body at  $\sim$ 1 AU in HD 69830, we must also posit a mechanism for moving the body inward after formation, e.g., planetary migration (Alibert et al. 2006). This mechanism may also be responsible for moving the Neptune-sized bodies inward as well. Analogous objects are the main belt asteroids in the 3:2 and 4:3 resonances with Jupiter. Another possibility is an orbital perturbation large enough to cause an icy parent body asteroid to be captured into the 2:1 or 5:2 HD 69830d resonances. Analogous objects are the Jupiter-family comets in the 1:2 resonances with Jupiter. The possibility of other, smaller planets in the system acting as efficient scatters of icy bodies in the outer reaches of HD 69830 is high, given the presence of three Neptune-sized planets, but until said planets are located, the size and nature of their scattering cannot be estimated.

# 6. CONCLUSIONS

Starting from a knowledge of small body mineralogy derived from the recent Deep Impact experiment excavation of comet 9P/Tempel 1, we are able to easily fit the observed disk excess emission observed by Beichman et al. (2005) using the Spitzer IRS. From this fitting, we find the following:

1. The IR excess is due to material in a dense disk orbiting around the star at 1.04  $\pm$  0.12 AU, coincident with the 2 :1 and

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5:2 orbital resonance with the outermost Neptune-sized planet, HD 69830d, and other stable low-eccentricity orbits with respect to the entire system.

2. The dust is totally lacking in amorphous pyroxenes, PAHs, and metallic sulfides and is deficient in Fe and S vs. solar abundances.

3. The material is experiencing ongoing collisional fragmentation, as evidenced by the continual supply of ephemeral small and icy dust particles.

4. The material originated in a parent P- or D-type asteroidal body originating most likely in an asteroid belt in the 2:1 resonance. The breakup of this body appears analogous to the fragmentation that accompanied the formation of the Karin and Veritas families in the solar system.

This paper was based on observations taken with the NASA Spitzer Space Telescope, operated by JPL/Caltech, and modeling performed under JPL contract 1274485. The authors would like to thank A. Cheng, O. Barnouin-Jha, C. Lovis, A. Rivkin, and A. Roberge for valuable input and R. Hurt for help with the figure graphics.

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