

# REDSHIFT SPACE DISTORTIONS : KAISER FORMULA AND BEYOND.

## LSS MINI COURSE

- INTRODUCTION
- KAISER FORMULA (SMALL ANGLE LINEAR REGIME)
- EXTENSION TO LARGE ANGLES
- NON LINEAR REGIME
- CONNECTING LINEAR AND NON LINEAR REGIMES
- OBSERVATIONS AND COSMOLOGICAL PARAMETERS
- CONNECTION WITH THE HALO MODEL

# INTRODUCTION

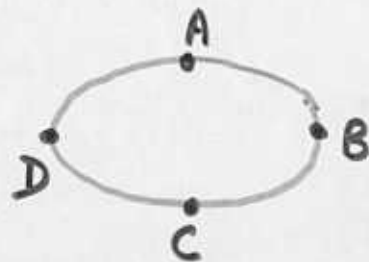
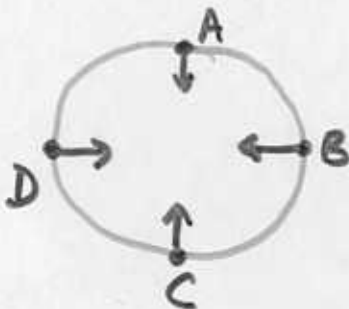
- REDSHIFT MEASUREMENT KEY TO LARGE SCALE STRUCTURE SURVEYS
- MEASURED REDSHIFT HAS TWO COMPONENTS
  - FROM THE HUBBLE EXPANSION
  - PECULIAR MOTION OF GALAXIES

$$v_{\text{MEASURED}} = \underbrace{Hd}_{\text{REAL SPACE EXPANSION}} + \underbrace{v_{\text{pec}}}_{\text{INFO. ABOUT } \phi\text{-POTENTIAL WELLS}}$$

REAL SPACE EXPANSION      INFO. ABOUT  $\phi$ -POTENTIAL WELLS

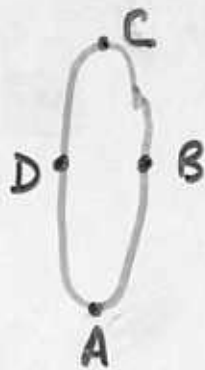
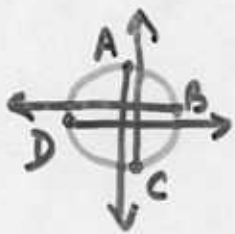
## — TWO EFFECTS

- LINEAR REGIME  $v_{\text{pec}} < Hd$



$$\begin{array}{c} \downarrow \\ v_{\text{pec}} \end{array} + \begin{array}{c} \uparrow \\ Hd \end{array} = \begin{array}{c} \uparrow \end{array} \quad : \quad \begin{array}{c} \uparrow \\ v_{\text{pec}} \end{array} + \begin{array}{c} \uparrow \\ Hd \end{array} = \begin{array}{c} \uparrow \end{array}$$

# - NON LINEAR REGIME



$$\downarrow + \uparrow = \downarrow$$

$v_{pec}$        $Hd$

$$\uparrow + \uparrow = \uparrow$$

$v_{pec}$        $Hd$

## KAISER FORMULA

— LINE OF SIGHT COMPONENT OF THE PECULIAR VELOCITY,  $\vec{u}(\vec{r})$

$$v = \vec{u}(\vec{r}) \cdot \hat{r}$$

— POSITION IN REDSHIFT SPACE,

$$s = r + v$$

$$\frac{ds}{dr} = 1 + \frac{\partial v}{\partial r}$$

— NUMBER OF GALAXIES IS CONSERVED,

$$n_s(s) d^3s = n_r(r) d^3r$$

$$\Rightarrow \bar{n}_s (1 + \delta^s(s)) s^2 ds = \bar{n}_r(r) (1 + \delta(r)) r^2 dr$$

$$\Rightarrow [1 + \delta^s(s)] = \frac{\bar{n}_r(r)}{\bar{n}_s(s)} \frac{1}{\left[1 + \frac{v}{r}\right]^2} \frac{1}{\left[1 + \frac{\partial v}{\partial r}\right]} [1 + \delta(r)]$$

—  $\bar{n}_r(r) = \bar{n}_s(s)$  AND TO LINEAR ORDER,

$$[1 + \delta^s(s)] \simeq [1 + \delta(r)] \left(1 - \frac{\partial v}{\partial r} - \frac{2v}{r} \dots\right)$$

— PLANE PARALLEL APPROXIMATION

$$\frac{\partial v}{\partial r} \gg \frac{2v}{r}$$

$$\Rightarrow \delta^s(s) \simeq \delta(r) - \frac{\partial v}{\partial r} \quad \text{--- ①}$$

ALSO IN LINEAR THEORY,

$$\vec{\nabla} \cdot \vec{u} + \frac{d\delta}{dt} = 0$$

$$\Rightarrow i\vec{k} \cdot \vec{u} + \delta \frac{d \ln \delta}{d \ln a} \frac{d \ln a}{dt} = 0$$

$$\Rightarrow \vec{u} = \frac{i\vec{k}}{k^2} \delta \frac{d \ln \delta}{d \ln a} \cdot \frac{d \ln a}{dt}$$

DEFINE  $\mu = \hat{r} \cdot \hat{k}$

$$\Rightarrow v = \hat{r} \cdot \vec{u} = \frac{i\mu}{k} \delta \underbrace{\frac{d \ln \delta}{d \ln a}}_{\text{GROWTH FACTOR}} \cdot \underbrace{\frac{d \ln a}{dt}}_{H=1}$$

$f(\Omega) \approx \Omega^{0.6}$

$$\Rightarrow v = -\frac{\mu}{k i} \delta f(\Omega)$$

$$\Rightarrow \frac{\partial v}{\partial r} = -\mu^2 \delta f(\Omega) \quad - (2)$$

USE (1) & (2)

$$\Rightarrow \delta^s(s) = \delta_l(r) [1 + \mu^2 f(\Omega)]$$

— PUT IN THE BIAS ASSUMING IT TO BE  
SCALE AND REDSHIFT INVARIANT

$$P_s(k) = [1 + \mu^2 \beta]^2 P(k)$$

— EXPAND IN LEGENDRE POLYNOMIALS

$$P_s(k) = \left[ \left( 1 + \frac{2}{3} \beta + \frac{1}{3} \beta^2 \right) P_0(\mu) + \left( \frac{4}{3} \beta + \frac{4}{7} \beta^2 \right) P_2(\mu) + \frac{8}{35} \beta^2 P_4(\mu) \right] P(k)$$

— KAISER FORMULA TURNS OUT TO BE  
GOOD ON VERY LARGE SCALES ( $\sim$  FEW 100 Mpc)  
SO NEED EXTENSIONS TO WORK WITH  
LSS SURVEY DATA SETS.

# EXTENSION TO LARGE ANGLES

(HEAVENS & TAYLOR '95; SZALAY, MATSUBARA & LANDY '98)

- SMALL ANGLE APPROACH OF KAISER WORKS ONLY AT SMALL WAVELENGTHS WHERE STRUCTURES ARE NONLINEAR
- NOW WE CAN USE ALL GALAXIES IN THE SURVEY
- AT LARGE ANGLES WE ARE IN THE LINEAR REGIME
- OUTLINE OF THE METHOD
  - EXPAND REDSHIFT SPACE DENSITY FIELD IN SPHERICAL HARMONICS

$$\tilde{P}_{\ell m n}(s) = c_{\ell n} \int p(s) \underbrace{j_{\ell}(k_{\ell n} s)}_{\text{SPHERICAL BESSEL FNC'S}} \underbrace{Y_{\ell m}^*(\theta, \phi)}_{\text{SPHERICAL HARMONICS}} d^3s$$

- NUMBER CONSERVATION OF GALAXIES

$$\Rightarrow p(s) d^3s = p(\gamma) d^3\gamma$$

$$\Rightarrow \tilde{P}_{\ell m n} = c_{\ell n} \int p(\gamma) j_{\ell}(k_{\ell n} s) Y_{\ell m}^*(\theta, \phi) d^3\gamma \quad - (1)$$

$$j_{en}(k_{en}s) = j_e(k_{en}(r+v))$$

$$\approx j_e(k_{en}r) + v(r) \frac{d}{dr} j_e(k_{en}r)$$

TO LINEAR ORDER

ALSO IN LINEAR THEORY,

$$\vec{v} = -\nabla\psi$$

$$\vec{\nabla} \cdot \vec{v} = -\delta f(\Omega)$$

$$\Rightarrow \nabla^2 \psi = \delta f(\Omega)$$

$$v(r) = \hat{r} \cdot \vec{v}$$

$$= f(\Omega) \sum_{lmn} \frac{c_{en} \delta_{emn} j'_e(k_{en}r) Y_{em}(\theta, \phi)}{k_{en}} - \textcircled{2}$$

PLUG  $\textcircled{2}$  IN  $\textcircled{1}$



# NON LINEAR REGIME (PEEBLES '80)

- EXCESS PROBABILITY OF FINDING TWO GALAXIES AT  $s$  AND  $s'$  IN REDSHIFT SPACE

$$dP = \bar{n}^2 d^3s d^3s' [1 + \xi(r_{\perp}, r_{\parallel})] \quad \text{--- (1)}$$

ANISOTROPIC IN REDSHIFT SPACE

- MEAN VALUE OF THE RELATIVE PECULIAR VELOCITY

$$\vec{v} = -g(r) \vec{v} ; \quad g(r) \begin{matrix} \rightarrow 1 \text{ ON SMALL SCALES} \\ \rightarrow 0 \text{ ON LARGE SCALES} \end{matrix}$$

- DISTRIBUTION OF VELOCITIES AROUND MEAN (ISOTROPIC)

$$dP_i = F(v) dv$$

- IF  $y$  IS THE TRUE SEPARATION ALONG LINE OF SIGHT

$$r_{\parallel} = y[1 - g(r)] + v_{\parallel}$$

- CONJECTURE

$$dP = \bar{n}^2 d^3s d^3s' [1 + \xi(r)] F(v) \delta^D[r_{\parallel} - y - g(r)y - v_{\parallel}] dv_{\parallel} \quad \text{--- (2)}$$

- EQUATE R.H.S. OF (1) & (2) AND PERFORM  $dv_{\parallel}$  &  $dy$  INTEGRALS

$$\xi(\gamma_{\perp}, \gamma_{\parallel}) = \int_{-\infty}^{\infty} dy \xi(\gamma) F_{\nu}[\gamma_{\parallel} - y + g(\gamma)y]$$

- ASSUME SOME FORM FOR  $F(\nu)$  TO OBTAIN MAP
- FOR A GAUSSIAN DISTRIBUTION
$$\delta_s(k) = \delta(k) e^{-\frac{(k\sigma\mu)^2}{2}}$$
- EXPONENTIAL AND LORENTZIAN WORK WELL TOO

# CONNECTING LINEAR & NONLINEAR REGIMES

— NONLINEAR FORMULATION IGNORES  $\delta v$  COUPLING

— ALSO ASSUMES THAT THE VELOCITY DISPERSION IS INDEPENDENT OF POSITION

— THE REGIME BETWEEN THE TWO EXTREMES IS MOST INTERESTING OBSERVATIONALLY.

— DEFINE,

$$\vec{\eta} = \begin{pmatrix} \delta(\vec{x}) \\ \delta(\vec{x}') \\ v(\vec{x}) \\ v(\vec{x}') \end{pmatrix}$$

$$[1 + \epsilon(\gamma_1, \gamma_2)] = \int d^4 \vec{\eta} dy (1 + \delta)(1 + \delta') F_{\vec{\eta}}(\vec{\eta}) \delta^D(\gamma_2 - y - v' + v)$$

— CAN DERIVE KAISER'S FORMULA FROM HERE BUT NONLINEARITY PREVENTS EXACT SOLUTION.

FUTURE . . . . .

# OBSERVATIONS AND COSMOLOGICAL PARAMETERS

— USEFUL FOR EXTRACTING REAL SPACE MAPS OF GALAXIES AND COMPARISON WITH THEORIES.

— MEASURING  $\beta \equiv \frac{\Omega^{0.6}}{b}$

— RATIO OF REDSHIFT SPACE TO REAL SPACE ANGLE AVERAGED POWER SPECTRA

$$\frac{P_S(k)}{P(k)} = 1 + 2\beta + \frac{1}{5}\beta^2$$

— RATIO OF QUADRAPOLE TO MONOPOLE HARMONICS OF POWER SPECTRUM

$$P_S(k) = \left[ \left(1 + \frac{2}{3}\beta + \frac{1}{5}\beta^2\right) P_0(l) + \left(\frac{4}{3}\beta + \frac{4}{7}\beta^2\right) P_2(l) + \frac{8}{35}\beta^2 P_4(l) \right] P(k)$$

$$P_0^S = \left(1 + \frac{2}{3}\beta + \frac{1}{5}\beta^2\right) P(k)$$

$$P_2^S = \left(\frac{4}{3}\beta + \frac{4}{7}\beta^2\right) P(k)$$

$$P_4^S = \frac{8}{35}\beta^2 P(k)$$

$$\frac{P_2^S(k)}{P_0^S(k)} = \frac{\frac{4}{3}\beta + \frac{4}{7}\beta^2}{1 + \frac{2}{3}\beta + \frac{1}{5}\beta^2}$$

— ASSUMES BIAS IS SCALE & REDSHIFT INDEPENDENT

— SOME RESULTS

— 2dF  $\beta = 0.43 \pm 0.07$

— IRAS  $\beta = 0.69^{+0.21}_{-0.19}$  (IR)

— STROMLO-APM  $\beta = 0.30^{+0.17}_{-0.15}$  (OPTICAL)

— CAN LEARN ABOUT GALAXY FORMATION

## CONNECTION WITH THE HALO MODEL

- LARGE SCALE POWER DOMINATED BY INTERACTION OF HALOS —  $P_{LIN}(k)$
- SMALL SCALE POWER DOMINATED BY INTERHALO CORRELATIONS —  $P_{NL}(k)$
- IN REDSHIFT SPACE

$$P_S(k) = P_{LIN}(k)(1+f(z)\mu)^2 + P_{NL}(k)e^{-(k\sigma\mu)^2/2}$$

- POWER AMPLIFIED ON LARGE SCALES DUE TO GRAVITATIONAL INFALL AND REDUCED ON SMALL SCALES DUE TO VIRIAL MOTIONS

## REFERENCES

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— OBSERVATIONS  
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